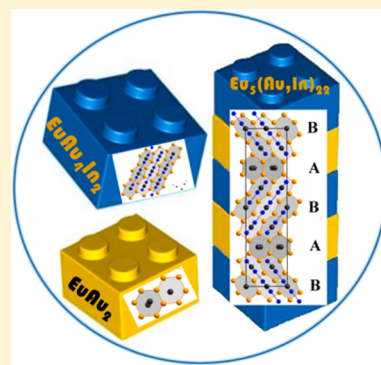


From the Ternary $\text{Eu}(\text{Au}/\text{In})_2$ and $\text{EuAu}_4(\text{Au}/\text{In})_2$ with Remarkable Au/In Distributions to a New Structure Type: The Gold-Rich $\text{Eu}_5\text{Au}_{16}(\text{Au}/\text{In})_6$ StructureSimon Steinberg,^{†,‡} Nathan Card,[‡] and Anja-Verena Mudring^{*,†,‡}[†]Ames Laboratory, U.S. Department of Energy and [‡]Department of Material Sciences and Engineering, Iowa State University, Ames, Iowa 50011, United States

S Supporting Information

ABSTRACT: The ternary $\text{Eu}(\text{Au}/\text{In})_2$ ($\text{EuAu}_{0.46}\text{In}_{1.54(2)}$) (I), $\text{EuAu}_4(\text{Au}/\text{In})_2$ ($\text{EuAu}_{4+x}\text{In}_{2-x}$ with $x = 0.75(2)$ (II), $0.93(2)$, and $1.03(2)$), and $\text{Eu}_5\text{Au}_{16}(\text{Au}/\text{In})_6$ ($\text{Eu}_5\text{Au}_{17.29}\text{In}_{4.71(3)}$) (III) have been synthesized, and their structures were characterized by single-crystal X-ray diffraction. I and II crystallize with the CeCu_2 -type (Pearson Symbol $oI12$; $Imma$; $Z = 4$; $a = 4.9018(4)$ Å; $b = 7.8237(5)$ Å; $c = 8.4457(5)$ Å) and the YbAl_4Mo_2 -type ($tI14$; $I4/mmm$; $Z = 2$; $a = 7.1612(7)$ Å; $c = 5.5268(7)$ Å) and exhibit significant Au/In disorder. I is composed of an Au/In-mixed diamond-related host lattice encapsulating Eu atoms, while the structure of II features ribbons of distorted, squared Au_8 prisms enclosing Eu, Au, and In atoms. Combination of these structural motifs leads to a new structure type as observed for $\text{Eu}_5\text{Au}_{16}(\text{Au}/\text{In})_6$ ($\text{Eu}_5\text{Au}_{17.29}\text{In}_{4.71(3)}$) ($oS108$; $Cmcm$; $Z = 4$; $a = 7.2283(4)$ Å; $b = 9.0499(6)$ Å; $c = 34.619(2)$ Å), which formally represents a one-dimensional intergrowth of the series EuAu_2 –“ EuAu_4In_2 ”. The site preferences of the disordered Au/In positions in II were investigated for different hypothetical “ $\text{EuAu}_4(\text{Au}/\text{In})_2$ ” models using the projector-augmented wave method and indicate that these structures attempt to optimize the frequencies of the heteroatomic Au–In contacts. A chemical bonding analysis on two “ EuAu_5In ” and “ EuAu_4In_2 ” models employed the TB-LMTO-ASA method and reveals that the subtle interplay between the local atomic environments and the bond energies determines the structural and site preferences for these systems.



INTRODUCTION

Polar metal-rich phases with gold have created excitement among scientists because of their unusual bonding patterns and rich, amazing structural chemistry, most recently the remarkable phases in the A–Au–Tr/Tt systems (A = alkaline/alkaline-earth; Tr/Tt = group 13 or 14 element).¹ The bonding patterns of these structures exhibit substantial $6s$ – $5d$ orbital mixings of gold as a consequence of relativistic effects,^{1–3} which have been observed for diverse more electron-poor intermetallics^{4–6} as well as for quasicrystals (QCs) and their approximants (ACs).^{7–11} Valence electron concentrations (vecs) and analyses of the band structures for these compounds utilizing density functional theory (DFT)-based methods place these structures between the Zintl and (close to) the Hume–Rothery phases.⁵ Typically, these more electron-poor intermetallic structures feature substantial Au/Tr or Au/Tt site mixing to optimize the bonding for the polyanionic networks, but still, the homogeneity ranges for these structures indicate certain electronic flexibilities.¹²

One class of compounds is the broad family of structures with diamond-like gold networks, which are observed for five different structure types: the rhombohedral $\text{Sr}_2\text{Au}_{6+x}\text{Zn}_{3-x}$ -type ($R\bar{3}c$),¹³ the hexagonal $\text{SrAu}_{4+x}\text{Al}_{3-x}$ -type ($P\bar{6}2m$),¹⁴ two independent orthorhombic structures, i.e. $\text{SrAu}_{5+x}\text{Al}_{2-x}$ ¹⁴ and BaAu_5Ga_2 -type ($Pnma$, both structures),¹⁵ and the monoclinic $\text{Sr}_2\text{Au}_{6+x}\text{Zn}_{3-x}$ -type ($C2/c$).¹³ The voids within

the Au host frameworks encompass triangular $[\text{Au},\text{M}]_3$ clusters ($\text{M} = \text{Al}, \text{Zn}, \text{Ga}, \text{In}, \text{or Sn}$) or alkaline-earth Ae (Sr, Ba) and Eu atoms,^{13–18} which are mutually exchanged for particular $\text{SrAu}_{4+x}\text{Al}_{3-x}$ -type structures. All these structures can be derived from the binary AAu_2 (A = Sr, Ba, Eu)^{19–21} to the extent that they are components of the Fibonacci series $(\text{Au}_{12/6}\text{A})_m(\text{Au}_{12/6}\text{M}_3)_n$ (A = Sr, Ba, Eu) with AAu_2 for $m = 1$, $n = 0$, AAu_4M_3 for $m = 1$, $n = 1$, $\text{A}_2\text{Au}_6\text{M}_3$ for $m = 2$, $n = 1$, and $\text{Ba}_{1.04}\text{Au}_{4.5}\text{Ga}_{2.4}$ or $\text{Eu}_{1.1}\text{Au}_{4.4}\text{Ga}_{2.2}$ for $m = 3$ and $n = 2$.¹⁵

Another class of Au-rich intermetallic compounds which show certain electronic flexibilities comprises those structures adopting the YbAl_4Mo_2 -type.^{22–26} This type of structure is composed of an aluminum network encapsulating ytterbium and molybdenum atoms, respectively.²⁴ While previous examinations concentrated primarily on the A–Au–Tr/Tt systems, the Au-rich parts of the R–Au–Tr/Tt phase diagrams (R = rare-earth element) have been studied to a lesser extent.^{27–29} Initial hints of rare-earth-containing $\text{SrAu}_{4+x}\text{Al}_{3-x}$ ¹⁵ and $\text{Sr}_2\text{Au}_{6+x}\text{Zn}_{3-x}$ -type¹⁸ compounds draw our interest to unprecedented realms of Au-rich, polar intermetallic compounds in the R–Au–Tr/Tt phase diagrams. Despite the fact that more recent research on the Eu–Au–In system has focused on its In-rich region,^{30–32} still, there are

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Table 1. Details of the Crystal Structure Investigations and Refinements for $\text{EuAu}_{0.5}\text{In}_{1.5}$, $\text{EuAu}_{4.8}\text{In}_{1.2}$, and $\text{Eu}_5\text{Au}_{17.3}\text{In}_{4.7}$

	$\text{EuAu}_{0.46}\text{In}_{1.54(2)}$ (I)	$\text{EuAu}_{4.75}\text{In}_{1.25(2)}$ (II)	$\text{Eu}_5\text{Au}_{17.29}\text{In}_{4.71(3)}$ (III)
fw	419.80	1231.08	4705.95
space group	<i>Imma</i> (No. 74)	<i>I4/mmm</i> (No. 139)	<i>Cmcm</i> (No. 63)
<i>a</i> , Å	4.9018(4)	7.1612(7)	7.2283(4)
<i>b</i> , Å	7.8237(5)		9.0499(6)
<i>c</i> , Å	8.4457(5)	5.5268(7)	34.619(2)
vol., Å ³	323.89(4)	283.43(7)	2264.6(2)
<i>Z</i>	4	2	4
density (calcd), g/cm ³	8.609	14.425	13.803
μ , mm ⁻¹	50.65	138.17	129.74
<i>F</i> (000)	700	999	7646
θ range	3.6–33.0	4.7–32.7	5.6–28.0
index ranges	$-8 \leq h \leq 9$ $-14 \leq k \leq 14$ $-14 \leq l \leq 15$	$-6 \leq h \leq 12$ $-11 \leq k \leq 11$ $-9 \leq l \leq 9$	$-8 \leq h \leq 8$ $-11 \leq k \leq 11$ $-42 \leq l \leq 42$
no. of reflns collected	3617	2054	10702
no. of independent reflns/ <i>R</i> _{int}	596/0.066	238/0.068	1233/0.154
no. of reflns with $I > 2\sigma(I)/R_\sigma$	456/0.043	196/0.035	792/0.084
refinement method		full-matrix least-squares on <i>F</i> ²	
data/restraints/parameter	596/0/13	238/0/11	1233/0/77
goodness-of-fit on <i>F</i> ²	1.04	1.11	1.01
final <i>R</i> indices [<i>F</i> ² > 2σ(<i>F</i> ²)]	<i>R</i> ₁ = 0.035; <i>wR</i> ₂ = 0.069	<i>R</i> ₁ = 0.026; <i>wR</i> ₂ = 0.049	<i>R</i> ₁ = 0.044; <i>wR</i> ₂ = 0.075
<i>R</i> indices (all data)	<i>R</i> ₁ = 0.053; <i>wR</i> ₂ = 0.077	<i>R</i> ₁ = 0.035; <i>wR</i> ₂ = 0.051	<i>R</i> ₁ = 0.086; <i>wR</i> ₂ = 0.090
largest diff. peak and hole, e ⁻ /Å ³	5.12 and -4.16	2.55 and -2.30	4.59 and -3.89

barely any reports regarding its Au-rich (more electron-poor) part.³³ Therefore, we present and discuss the results of the initial explorations for the Au-rich part of the Eu–Au–In system, which yielded the intermetallic compounds $\text{EuAu}_{0.5}\text{In}_{1.5}$, $\text{EuAu}_{4+x}\text{In}_{2-x}$ ($x = 0.8–1.0$), and $\text{Eu}_5\text{Au}_{17.3}\text{In}_{4.7}$ with a new structure type for the last composition.

EXPERIMENTAL TECHNIQUES

Syntheses. Gold pieces (99.999%, BASF), filings of europium (99.99%, Ames Laboratory), indium and gallium ingots (99.999% Alfa Aesar) were used as starting materials and stored and handled under dry argon atmosphere in a glovebox ($\text{H}_2\text{O} < 0.1$ ppmv). The Eu filings were produced from large ingots, which were mechanically polished prior to each use. Loads of ~200–300 mg total were weighed in pre-cleaned, one-side arc-welded tantalum tubes, which were closed inside a glovebox, arc welded on the other end, and jacketed by evacuated Schlenk flasks. The products were obtained using the following temperature programs. $\text{EuAu}_{0.5}\text{In}_{1.5}$: heat to 800 °C within 4 h, keep that temperature for 12 h, cool to 300 °C with 10 °C/h, and anneal at this temperature for 60 h. $\text{EuAu}_{4+x}\text{In}_{2-x}$ ($x = 0.8$) and $\text{Eu}_5\text{Au}_{17.3}\text{In}_{4.7}$: heat to 800 °C in 4 h, keep the temperature for 12 h, slowly cool to 300 °C (2 °C/h), anneal at this temperature for 3 days, and quench in water. Single crystals obtained from samples with loads corresponding to hypothetical compositions of “ $\text{Eu}_2\text{Au}_7\text{In}_2$ ” and “ $\text{Eu}_1\text{Au}_5\text{In}_1$ ” (Supporting Information; Table S1) pointed to a narrow solid solution for $\text{EuAu}_{4+x}\text{In}_{2-x}$ ($x = 0.8–1.0$). The products appeared as gray powders with silver crystals of polyhedral shapes and metallic luster and were stable in air in accord with recently reported Au-rich compounds.^{15,16} Analogous reactions in the Eu–Au–Ga system yielded a rhombohedral $\text{Sr}_2\text{Au}_{6+x}\text{Zn}_{3-x}$ -type¹³ phase ($\text{Eu}_2\text{Au}_{6.1}\text{Ga}_{2.9}$; *R*-3c; *a* = 8.411(5) Å; *c* = 21.830(13) Å, see Supporting Information) or hexagonal $\text{SrAu}_{4+x}\text{Al}_{3-x}$ -type¹⁴ compounds, which have been reported elsewhere.¹⁵

X-ray Studies. All samples were checked for purity through detailed phase analyses of sets of powder X-ray diffraction data, which were collected on STOE STADI P diffractometers equipped with a STOE image plate and DECTRIS MYTHEN 1K detectors (Cu *K*_{α1}; $\lambda = 1.54059$ Å) at room temperature. For the measurement the samples were dispersed on Mylar sheets with grease and placed between split

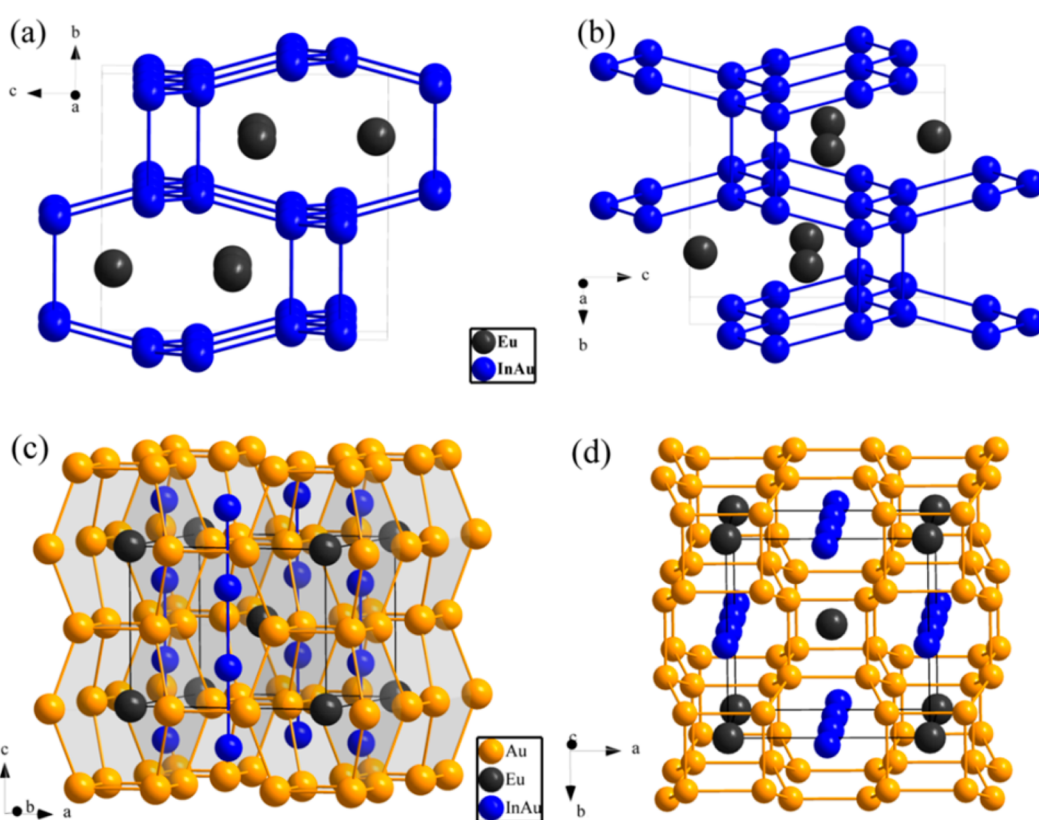
aluminum rings or sealed in capillaries. Data and phase analyses were carried out with the *WinXPow* software package.³⁴ Detailed phase analyses of the powder patterns revealed that $\text{EuAu}_{4+x}\text{In}_{2-x}$ ($x = 0.75(2)$ (II), $0.93(2)$, and $1.03(2)$) were obtained as high-yield products in competition with $\text{Eu}_5\text{Au}_{17.3}\text{In}_{4.7}$ (III), while $\text{EuAu}_{0.5}\text{In}_{1.5}$ (I) was accompanied by the previously reported EuAuIn ³⁵ (see Supporting Information).

Single crystals were selected from the bulk samples, fixed on glass fibers with grease, and subsequently transferred to a Bruker APEX CCD diffractometer (Bruker, Inc.; Madison, WI). Sets of single-crystal X-ray intensity data were collected in φ - and ω -scan modes at room temperature (~296 K) with Mo *K* α radiation ($\lambda = 0.71073$ Å) and exposures of 15 s/frame. The raw intensity data were integrated utilizing the *SAINTE* program within the *SMART* software suite,³⁶ while the program *SADABS*³⁷ was employed for empirical absorption corrections.

Checking the intensity data sets for extinction conditions with the *XPREP* algorithms in the *SHELXTL* suite³⁸ and examination of the *E*²-*I* statistics lead to the assignment of the centrosymmetric space groups *Imma* (No. 74), *I4/mmm* (No. 139), and *Cmcm* (No. 63) for $\text{EuAu}_{0.46}\text{In}_{1.54(2)}$ (I), $\text{EuAu}_{4+x}\text{In}_{2-x}$ ($x = 0.75(2)$ (II), $0.93(2)$, and $1.03(2)$), and $\text{Eu}_5\text{Au}_{17.29}\text{In}_{4.71(3)}$ (III), respectively. All structures were solved using direct methods (*SHELXS-2013*) and refined on *F*² in full-matrix least-squares including anisotropic atomic displacement parameters (*SHELXL-2013*).³⁹ Initial refinements of the *8h* sites in I and the *4d* positions in II with the scattering factor of indium lead to too small displacement parameters for these sites, and close inspection of the Fourier maps for I and II indicated higher electron densities on these sites; however, subsequent refinements with the scattering factors of gold resulted in too large anisotropic atomic displacement parameters for the *8h* sites in I and the *4d* positions in II. Accordingly, mixed Au/In occupancies were introduced for these sites, and approaches to refine them with the same as well as different positional parameters for both gold and indium preceded to more reasonable anisotropic atomic displacement parameters for the refinements with the common Au/In positions. Analogous examinations of the *16h* and *8f* positions for the orthorhombic structure of III indicated the presence of mixed Au/In occupancies for the latter sites. Therefore, mixed Au/In occupancies were assigned on three *8f* positions in the structure of III, which was refined to the composition

Table 2. Atomic Positions and Equivalent Isotropic Displacement Parameters for $\text{EuAu}_{0.46}\text{In}_{1.54(2)}$, $\text{EuAu}_{4.75}\text{In}_{1.25(2)}$, and $\text{Eu}_3\text{Au}_{17.29}\text{In}_{4.71(3)}$

atom	position	<i>x</i>	<i>y</i>	<i>z</i>	$U_{\text{iso}}/U_{\text{eq}}$ Å ²	Occ. (<1)
$\text{EuAu}_{0.46}\text{In}_{1.54(2)}$						
Eu1	4e	0	1/4	0.5418(1)	0.0170(2)	
Au2/In2	8h	0	0.0485(1)	0.1637(1)	0.0181(2)	0.23(1)/0.77(1)
$\text{EuAu}_{4.75}\text{In}_{1.25(2)}$						
Au1	8h	0.3059(1)	0.3059(1)	0	0.0129(2)	
Eu2	2a	0	0	0	0.0112(3)	
Au3/In3	4d	0	1/2	1/4	0.0156(4)	0.37(1)/0.63(1)
$\text{Eu}_3\text{Au}_{17.29}\text{In}_{4.71(3)}$						
Au1	16h	0.1933(2)	0.3090(1)	0.4620(1)	0.0108(3)	
Au2	16h	0.1975(2)	0.4228(1)	0.2901(1)	0.0089(3)	
Au3	16h	0.1931(2)	0.0630(1)	0.4136(1)	0.0108(3)	
Au4	16h	0.3064(2)	0.1765(1)	0.3367(1)	0.0115(3)	
Au5/In5	8f	0	0.4057(3)	0.5319(1)	0.0085(10)	0.20(2)/0.80(2)
Au6/In6	8f	1/2	0.2846(3)	0.4045(1)	0.0119(11)	0.37(2)/0.63(2)
Au7/In7	8f	1/2	0.4833(4)	0.3395(1)	0.0165(14)	0.07(2)/0.93(2)
Eu8	4a	1/2	1/2	1/2	0.0091(7)	
Eu9	8f	0	0.3703(3)	0.3755(1)	0.0093(6)	
Eu10	4c	0	0.7087(4)	1/4	0.0065(7)	
Eu11	4c	0	0.1472(4)	1/4	0.0101(8)	

**Figure 1.** Representation (a) of the unit cell of I: Au and In atoms form hexagons in chair conformations within the *ac* plane (b). (c) View on the unit cell of II: each europium atom is enclosed by 12 Au and 8 M atoms, while the gold network may be depicted as mutually perpendicular gold hexagons in boot conformations (d).

$\text{Eu}_3\text{Au}_{17.29}\text{In}_{4.71(3)}$ for the selected single crystal. The *PLATON* software package⁴⁰ was employed to check the symmetries of the refined structures, and no higher symmetries were identified. With these settings the R_1 values converged to less than 4.36% (Table 1). Details of the data collection and refinement parameters for I, II, and III are provided in Table 1, whereas atomic positions are listed in Table 2.

Computational Details. Electronic structure calculations were undertaken for hypothetical $\text{EuAu}_4(\text{Au}/\text{In})_2$ (II) models representing the compositions “ EuAu_4In_2 ” and “ EuAu_3In ” (Figure 3) to identify the site preferences in the tetragonal structures. The “ EuAu_4In_2 ” model differs from the experimentally determined tetragonal $\text{EuAu}_{4+x}\text{In}_{2-x}$ structures ($x = 0.8-1.0$) through full indium occupation of the *4d* site, while in the “ EuAu_3In ” models indium atoms reside on both the *8h* and *4d* sites. Full structural optimizations of all models and band

structure calculations were carried out using the projector-augmented wave (PAW) method of Blöchl⁴¹ as implanted in the *Vienna ab initio Simulation Package* (VASP) by Kresse and Joubert.^{42–46} Correlation and exchange were described by the general gradient approximation of Perdew, Burke, and Enzerhof (GGA–PBE),⁴⁷ while starting meshes of $4 \times 4 \times 4$ up to $8 \times 8 \times 8$ k points were employed to sample the first Brillouin zone for reciprocal space integrations. The cutoff energy of the plane wave basis sets was set 500 eV, and full optimizations were accomplished until the energy difference between two iterative steps fell below 10^{-7} eV/cell. An effective on-site Coulomb interaction term ($U_{\text{eff}} = 3$ eV) was included in the Kohn–Sham Hamiltonian to account for the strong correlations within the Eu $4f$ states.

A chemical bonding analysis on all models was completed based on the integrated values of the crystal orbital Hamilton populations (COHP), which were obtained through the tight-binding linear-muffin-tin-orbital (TB-LMTO) method with the atomic sphere approximation (ASA) using the Stuttgart code.^{48,49} In particular, the off-site projected density of states are weighted by the corresponding Hamilton matrix elements to determine bonding and antibonding states.⁵⁰ The optimized structural parameters of the models were obtained from the VASP computations and used as starting points for the TB-LMTO-based electronic structure calculations. The Wigner–Seitz (WS) radii were generated automatically, and empty spheres (ES) were included to achieve an optimal approximation of full potentials. The basis set employed the following orbitals (downfolded⁵¹ orbitals in parentheses): Au-6s/-6p/-5d/(-5f); Eu-6s/(-6p)/-5d; In-5s/-5p/(-5d)/(-4f). The corresponding WS radii (Angstroms) were as follows: Au, 2.91–3.13; Eu 3.49–3.72; In, 3.02–3.14. On the basis of the VASP computations the Eu $4f$ states are strongly localized and, hence, treated as core-like states in the TB-LMTO-based calculations, which is in accord with recently reported results on R-containing compounds.^{52–54} Reciprocal space integrations were completed with the tetrahedron method⁵⁵ employing $24 \times 24 \times 24$, $24 \times 24 \times 24$, $24 \times 24 \times 24$, $12 \times 6 \times 12$, and $12 \times 12 \times 6$ k -point sets for the “EuAu₄In₂”, “EuAu₃In-1”, “EuAu₃In-2”, “EuAu₃In-3”, and “EuAu₃In-4” models, respectively. Plots of the DOS and –COHP curves are shown in Figure 4, while the integrated values of the –COHP curves (ICOHP) may be extracted from the Supporting Information (Tables S3–S6).

RESULTS AND DISCUSSION

More recent investigations of the Au-rich part of the Eu–Au–Tr (Tr = Al, Ga) systems have resulted in the identification of rhombohedral Sr₂Au_{6+x}Zn_{3-x}-type^{13,18} (Table S1, Supporting Information) or hexagonal SrAu_{4+x}Al_{3-x}-type^{14,15} compounds. Initial examinations of the Eu-poor regions (<33 at.%) for the corresponding Eu–Au–In system lead to the discovery of the Au-rich EuAu_{4+x}In_{2-x} ($x = 0.75(2)$ (II)) and Eu₅Au_{17.29}In_{4.71(3)} (III), which were detected at the same composition “Eu₂Au₇In₂”. Furthermore, the ternary EuAu_{0.46}In_{1.54(2)} (I) was also found at this composition. Subsequent reactions of samples loaded as “EuAu₃In” produced high yields of EuAu_{4+x}In_{2-x} ($x = 0.8–1.0$; Figure S3, Supporting Information), while loadings of “EuAu₄In₂” uncovered the possibility of Tsai-type quasicrystals and their approximants in this system.

Crystal Structures. The structure of EuAu_{0.46}In_{1.54(2)} (I) is derived from the binary EuAu₂ (*Imma*; No. 74; CeCu₂-type⁵⁶) through partial indium substitution on the gold sites (Wyckoff position $8h$).²¹ In particular, this structure features puckered sheets of gold/indium-mixed (M) hexagons adopting chairlike conformations within the *ac* plane. The M–M distances are short (2.851(1)–2.868(1) Å) within the M₆ rings and range from 3.152(1) to 4.671(1) Å between the M hexagons sandwiching the Eu host atoms (Wyckoff position $4e$). Accordingly, each europium atom is enclosed by 12 M atoms forming condensed Eu@M_{12/6} cages (Figure 1). Note that the structure of I crystallizes with the CeCu₂-⁵⁶ rather than the

CaIn₂-type (P6₃/*mmc*; No. 194), which is adopted by the binary EuIn₂.⁵⁷ Alike Au/In distributions have also been observed for the polyanionic networks of the previously reported BaAu_{0.4}In_{1.6}⁵⁸ and GdAu_{0.4}In_{1.6}.⁵⁹ The structure of the barium-containing compound represents a substitutional solution of gold in BaIn₂⁶⁰ (CeCu₂-type⁵⁶), while the M network of GdAu_{0.4}In_{1.6} (CaIn₂-type⁵⁷) assembles a hexagonal diamond-like substructure as observed for the rhombohedral Eu₂Au_{6.1}Ga_{2.9} (see Supporting Information).

Alloys in the series EuAu_{4+x}In_{2-x} ($x = 0.8–1.0$) adopt the YbAl₄Mo₂-type of structure (*I4/mmm*; No. 139)²⁴ and are composed of gold networks encapsulating Eu and M (M = Au/In) atoms, respectively. The structures of the EuAu_{4+x}In_{2-x} ($x = 0.8–1.0$) series show alike structural features, and the indium-rich representative, EuAu_{4.75}In_{1.25(2)}, was selected for further structural analysis. Gold atoms residing on Wyckoff position $8h$ form mutually perpendicular, puckered layers of Au₆ hexagons, which adopt boot-like conformations within the *ac* and *bc* planes (Figure 1). The Au–Au distances along the apexes of the hexagons (2.986(1) Å) are longer than the Au–Au separations (2.780(1) Å) that encompass the bases of the Au₆ circles and, furthermore, the Au₄ squares parallel to the *ab* plane. Alternatively, the Au substructure can be depicted as distorted, squared Au₈ prisms that are condensed via common faces to construct ribbons along (001). Each europium atom (Wyckoff position $2a$) is enclosed by two face-sharing Au₈ prisms with eight M atoms (Wyckoff position $4d$) capping their faces along the *c* axis (Figure S5, Supporting Information). Accordingly, each europium atom is coordinated by 12 Au and eight M atoms, guiding to a coordination number of 20 for Eu. Similarly, each M atom occupies a center of a distorted, squared Au₈ prisms with Au–M distances of 2.939(1) Å. Additionally, the M@Au₈ cages are bicapped by the endohedral M atoms of the nearest neighboring Au₈ prisms leading to M–M contacts of 2.763(1) Å within the ribbons.

The nearest structural analogues come from the previously reported SrAu_{4.8}In_{1.2}²⁵ and EuAu_{4.8}Cd_{1.2}²⁶ (YbAl₄Mo₂-type;²⁴ both structures). The Sr-containing structure features close M–M contacts (M = Au/In; 2.764(1) Å) within the chains and Au–Au distances of 2.782(1) Å for the Au₄ squares, while the M–M distances (here M = Au/Cd) are 2.730(1) Å within the chains in the structure of EuAu_{4.8}Cd_{1.2}. A clear difference between SrAu_{4.8}In_{1.2} and the structures of EuAu_{4+x}In_{2-x} ($x = 0.93(2)$, 1.03(2); Supporting Information) arises from the distribution of the Au and In atoms, as the mixed $4d$ sites are occupied by about or even more than 50 atom % Au in the Eu-containing compounds. Comparatively, phase analytical studies based on powder X-ray diffraction data for the Eu–Au–Cd systems pointed to a substitutional homogeneity range of $0 \leq x \leq 1$ for EuAu_{4+x}Cd_{2-x},²⁶ yet, a “EuAu₄In₂” composition, which shows no Au/In mixing for the $4d$ sites, has not been observed for this system. Previous investigations in the Ce–Mn–Cu⁶¹ and La–Ag–Mg⁶² systems indicated that the ranges of compositions for solid solutions with the YbAl₄Mo₂ type of structure partially overlap with those of different structure types. For instance, an increase of the Ag content in LaAg_{4+x}Mg_{2-x} going from $x = 0.4$ to $x = 0.9$ leads to a different chemical occupation of the relative crystallographic sites and an order–disorder transformation from the YbAl₄Mo₂- (*I4/mmm*; No. 139) to the LaAg₅Mg-type (*P4/nmm*; No. 129).⁶² Such polymorphism could not be encountered for the structures of EuAu_{4+x}In_{2-x} ($x = 0.8–1.0$), as close inspection of the extinction conditions for the X-ray

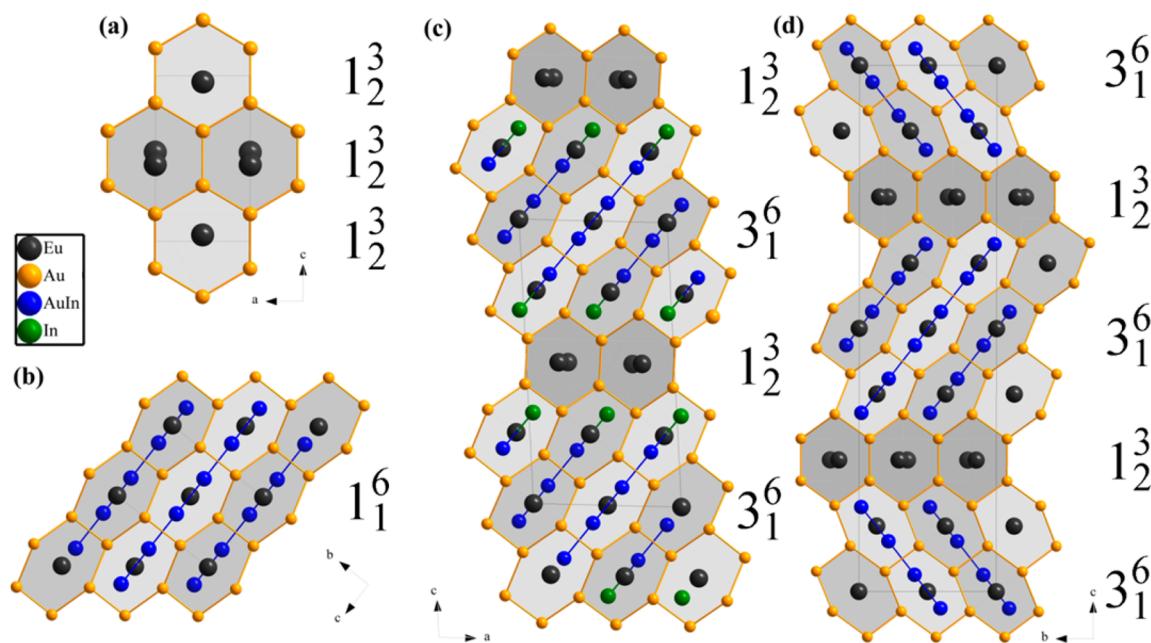


Figure 2. Representatives of the linear, inhomogeneous intergrowth series EuAu_2^{21} (1_2^3 ; a)–“ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ” (1_1^6 ; b; Eu, dark gray; Au, yellow; M, blue; In, green): the monoclinic $\text{Eu}_5\text{Au}_{17.7}\text{In}_{4.3}$ ($C2/m$; c) and orthorhombic $\text{Eu}_5\text{Au}_{17.3}\text{In}_{4.7}$ ($Cmcm$; d) formally represent one-dimensional intergrowths of the series EuAu_2 –“ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ”. Full indium occupancies of the disordered $4d$, $4i$, and $8f$ sites for $\text{EuAu}_{4.8}\text{In}_{1.2}$, $\text{Eu}_5\text{Au}_{17.7}\text{In}_{4.3}$, and $\text{Eu}_5\text{Au}_{17.3}\text{In}_{4.7}$, respectively, lead to the compositions “ EuAu_4In_2 ” and “ $\text{Eu}_5\text{Au}_{16}\text{In}_6$ ”, which point to the presence of an intergrowth series. The two-dimensional parent segments EuAu_2 and “ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ” form infinite slabs stacked along (001) for both intergrowths. Note the dissimilar stacking between the monoclinic (c) and the orthorhombic (d) structures.

intensity data sets as well as E^2-1 statistics for these structures clearly pointed to the space group $I4/mmm$. It is worth noting that mixed occupancies occur on the $4d$ rather than the $8h$ sites in the crystal structures of the YbAl_4Mo_2 -type ternary phases with substitutional homogeneity ranges;^{25,26,61–63} however, conclusive hints accounting for this preference have not been provided. Electronic structure calculations on different “ EuAu_5In ” models will provide insight into the origin for this site preference (see [Coloring and Electronic Structure](#)).

$\text{Eu}_5\text{Au}_{17.29}\text{In}_{4.71(3)}$ (III) crystallizes in its own structure type ($Cmcm$; No. 63) and formally resembles a one-dimensional (or linear) intergrowth of “ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ”-related (the term “ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ” refers to segments showing the similar structural features as observed for $\text{EuAu}_{4+x}\text{In}_{2-x}$ ($x = 0.8$ – 1.0)) and EuAu_2 -related (two-dimensional) slabs. Strictly speaking, evaluation of a direct structural relationship between $\text{EuAu}_{4.75}\text{In}_{1.25(2)}$ and $\text{Eu}_5\text{Au}_{17.29}\text{In}_{4.71(3)}$ is set hurdles by the dissimilar occupancies of the disordered positions in II and III; however, a topological analysis of all structures with an approximation of full indium occupancies for the mixed Au/In $4d$ and $8f$ positions in II and III, respectively, guides to a structural relationship akin to a linear inhomogeneous intergrowth from EuAu_2 and II to III.

The orthorhombic structure contains three independent M sites (Wyckoff positions $8f$, $8f$, $8f$) residing in the centers of the distorted, squared Au_8 prisms. The contacts between two endohedral M atoms ($2.790(4)$ – $2.879(4)$ Å) are slightly larger than those observed for II and suggest coordination environments of eight Au and two M atoms for the four M atoms within one ribbon (Figure S6, [Supporting Information](#)); however, the M@Au_8 cages at the ends of each ribbon are capped by one M and one Eu atom ($d(\text{Eu}–\text{M}) = 3.436(4)$ Å). In addition to the intercalation of the EuAu_2 -fashioned slabs, the orthorhombic structure comprises four independent

europium positions. Two europium sites (Wyckoff positions $4c$, $4c$) are located in the Au_{12} cages of the EuAu_2 segments, whereas two other europium sites (Wyckoff position $4a$ and $8f$) are enclosed by face-sharing Au_8 prisms (Figure S6, [Supporting Information](#)). The Eu atoms in the centers of the “ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ”-related slabs (Wyckoff position $4a$) are surrounded by 12 Au and eight M atoms in accord with the Eu coordination spheres in II, while the europium atoms residing on Wyckoff position $8f$ possess (12Au + 7M) coordination environments. In contrast to the tetragonal structure the bases of the Au_8 prisms show slight distortions, as they range from $2.785(2)$ to $2.794(2)$ Å within and from $2.799(2)$ to $2.856(2)$ Å at the edges of the “ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ” segments. The Au–Au contacts in and between the Au_6 rings of the EuAu_2 layers scale from $2.776(2)$ to $2.861(2)$ Å, in fair agreement with those observed for EuAu_2^{21} (2.714 – 2.917 Å).

The nearest structural analogue comes from the previously reported $\text{Eu}_5\text{Au}_{17.7}\text{In}_{4.3}$ ($C2/m$; No. 12),³³ which is also composed of EuAu_2 - and “ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ” parental segments. Although a strict structural comparison between II, III, and the monoclinic structure cannot be accomplished due to the dissimilar disorder of the M positions, a topological analysis of the segments for $\text{Eu}_5\text{Au}_{17.3}\text{In}_{4.7}$ and $\text{Eu}_5\text{Au}_{17.7}\text{In}_{4.3}$ implies that these structures resemble one-dimensional (or linear) intergrowths in the series EuAu_2 –“ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ” based on the approximations of full indium occupancies for the disordered M sites in these structures. In particular, the orthorhombic $\text{Eu}_5\text{Au}_{17.3}\text{In}_{4.7}$ and the monoclinic $\text{Eu}_5\text{Au}_{17.7}\text{In}_{4.3}$ both feature sequences of three succeeding “ $\text{EuAu}_{4+x}\text{In}_{2-x}$ ” segments (3_1^6) and one EuAu_2 slab (1_2^3), which are stacked alternately along (001) in accordance with the formulas $(3_1^6 1_2^3)$ and $(3_1^6 1_2^3)_2$ for the monoclinic and orthorhombic structures, respectively (Figure 2). The superscript indexes 3 and 6 represent mirror planes and two-folded axes, respectively, that are arranged

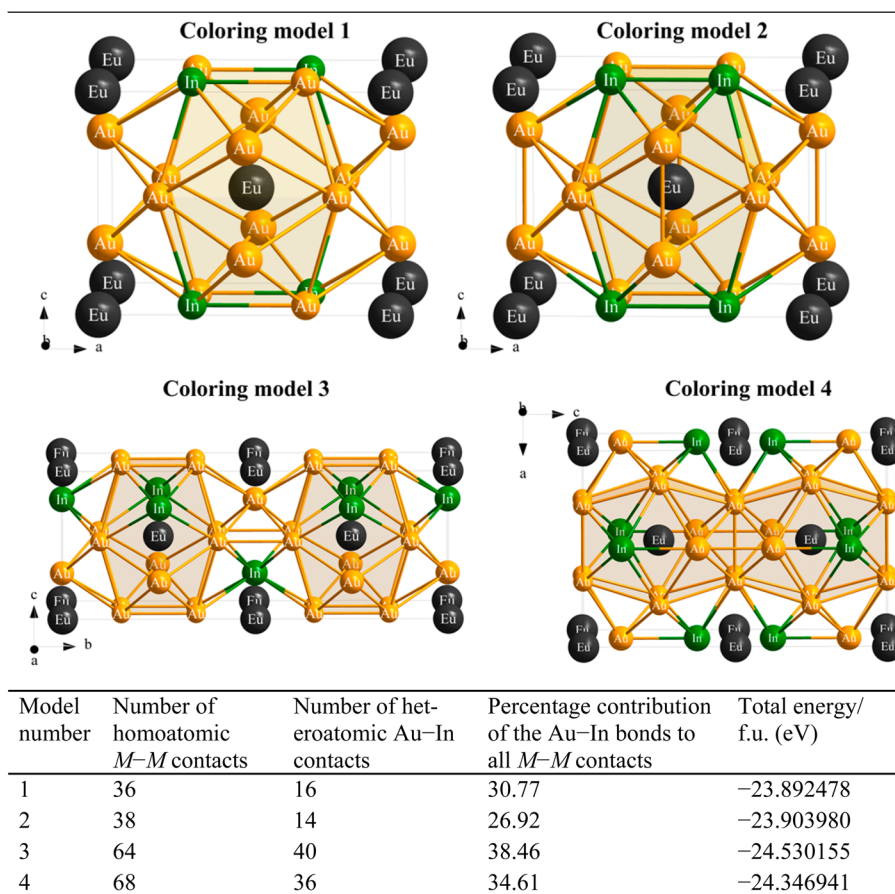


Figure 3. Coloring models of hypothetical “EuAu₅In” structures: gold and indium sites are represented by the gold and green atoms, respectively. In models 1 and 2 (top) the indium atoms reside on the 8*h* sites, while the 4*d* positions are occupied by indium in models 3 and 4 (bottom). From the total energies of the different schemes it is clear that there is a site preference for the 4*d* sites. Placing In on the 4*d* sites surrounded by 10 nearest neighboring gold atoms leads to the maximum possible number of 10 heteroatomic Au–In contacts per In. Because each model comprises 26 $M-M$ contacts/f.u., the maximal percentage of the Au–In bonds to all $M-M$ contacts is reached for 38.46%, which is achieved for “EuAu₅In-3”. Models 3 and 4 are constructed of two “EuAu_{4+x}In_{2-x}” segments to accomplish broader options for structural variations.

perpendicular to the stacking direction within one segment.⁶⁴ In the orthorhombic structure the neighboring layers of the three succeeding “EuAu_{4+x}In_{2-x}” parental segments (3_1^6) are aligned in opposite directions with a torsion angle of 102.7°, while the “EuAu_{4+x}In_{2-x}” slabs line up in the same direction in the monoclinic structure (Figure 2). Accordingly, the orthorhombic structure comprises two ($3_1^6 1_2^3$) sections in succession and has a longer translation period along the stacking direction than the monoclinic compound. The disordered M sites of the orthorhombic structure exhibit larger indium contributions relative to the monoclinic compound and lead to a higher indium content in III than for the monoclinic Eu₅Au_{17.7}In_{4.3}.

Coloring and Electronic Structure. An analysis of the electronic structure for the ribbons of the face-sharing Au₈ cages enclosing Eu, Au, and In atoms was accomplished for the tetragonal structures, which show the shortest repeat unit along (001) within the series EuAu_{4.8}In_{1.2}, Eu₅Au_{17.7}In_{4.3}, and Eu₅Au_{17.3}In_{4.7}. To understand the site preference for the mixed Au/In occupancy on the Wyckoff position 4*d* different starting models were examined (see Computational Details): an ordered “EuAu₄In₂” model that is derived from the tetragonal structures through full indium occupations of the mixed M sites and diverse “EuAu₅In” models with mixed Au/In occupations on both the 8*h* (models 1 and 2) and the 4*d* (models 3 and 4)

sites, respectively (Figure 3). The diverse “EuAu₅In” models were inspected to reveal the influence of dissimilar heteroatomic Au–In bond frequencies and different local atomic environments on the arrangements of gold and indium in the tetragonal structures. In intermetallic compounds with disordered networks the model with the maximal number of heteroatomic contacts tends to provide the lowest total energy and is typically expected to be the most favorable scheme to compare to the experimental structure.^{65–67} To examine a likely presence of this tendency for the tetragonal compounds, schemes with maximal numbers of heteroatomic Au–In contacts have been achieved for the gold and indium distributions in models 1 and 3, which resemble Au/In partitioning on the 8*h* and 4*d* sites of the observed structure, respectively. The Au/In distributions in “EuAu₅In” models 2 and 4 give slightly lower amounts of Au–In contacts relative to “EuAu₅In-1 and -3”, respectively, to develop tendencies between the bond frequencies, local atomic environments, and bond as well as total energies for this series. The outcome of this analysis will help to understand the site preferences for this particular system and to corroborate the experimentally determined structure as the electronically most favorable.

A topological analysis of the tetragonal structure reveals that an indium occupation of the 8*h* site brings about a maximum of 8 heteroatomic Au–In contacts per indium atom, while placing

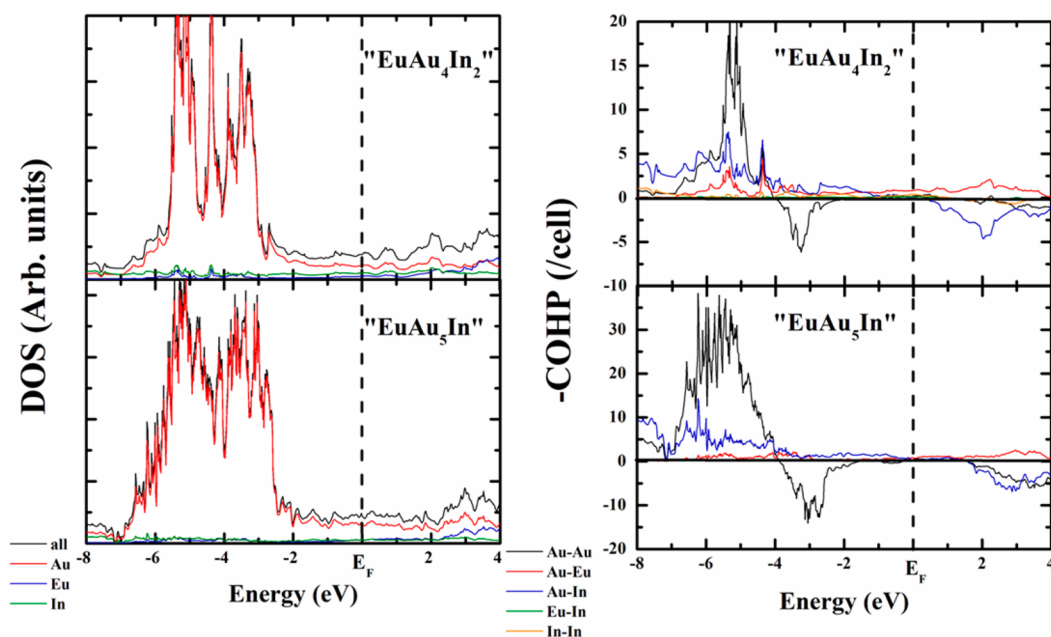


Figure 4. DOS (left) and $-\text{COHP}$ curves (right) of the “EuAu₄In₂” and “EuAu₅In” models computed by the TB-LMTO method: the Fermi level is represented by the dashed lines, while orbital-projected DOS curves of both models may be extracted from the [Supporting Information](#) (Figure S9). Optimized structural parameters of the “EuAu₅In” model 3, which is the “EuAu₅In” composition with the lowest total energy, were used as starting points for the TB-LMTO-based calculations. Spin-polarized DOS curves (VASP) are provided in Figure S8, [Supporting Information](#).

indium on the Wyckoff position $4d$ leads to a maximum of 10 Au–In contacts/In. Accordingly, higher Au–In bond frequencies are achieved for all M–M contacts in models 3 and 4, in which the indium atoms reside on the positions $4d$, rather than for models 1 and 2 with In atoms occupying the $8h$ sites (Figure 3). Because the scheme with the maximal number of heteroatomic contacts tends to give the lowest total energies for structures with disordered networks,⁶⁸ one would expect the model with the lowest homoatomic bond frequency to be the most favorable to compare to the experimental structure. A comparison of the total energies for the “EuAu₅In” models indicates that the lowest total energy is obtained for model 3, which is the scheme with the largest number of Au–In contacts per indium atom, and provides justification why indium prefers to occupy the $4d$ rather than the $8h$ sites.

Because preliminary calculations with varying correlation parameters (Figure S7, [Supporting Information](#)) indicated strong correlations for the Eu $4f$ states,^{69–71} an effective on-site Coulomb interaction term ($U_{\text{eff}} = 3$ eV) was included as a correctional parameter to the Kohn–Sham Hamiltonian of the spin-polarized calculations. The total energy calculations on a ferromagnetic and an antiferromagnetic “EuAu₄In₂” model revealed an energy difference of 8.2 meV/cell between the antiferromagnetic and the lower lying ferromagnetic state, which tends to be preferred in fair agreement with the magnetic data.⁷² The spin-polarized DOS curves of “EuAu₄In₂” (Figure S8, [Supporting Information](#)) do not superimpose, inferring a ferromagnetic state with a major difference for the Eu $4f$ states as the origin of a magnetic response.⁷² The bands accounting for the Eu $4f$ AOs exhibit extremely small dispersions leading to sharp peaks at 4.22 and 11.72 eV in the spin-up and spin-down DOS curves, respectively. Such small dispersions are indicative of rather localized states playing a subordinate role in overall bonding and provide justification to treat the Eu $4f$ AOs as core-like states in the LMTO-based calculations.

Density of States. To provide insight into the electronic structures of the disordered tetragonal compositions EuAu_{4+x}In_{2–x} ($x = 0.8–1.0$), which lie between those of the ordered “EuAu₄In₂” and “EuAu₅In”, we followed up with an analysis of the DOS curves (Figure 4) for the “EuAu₄In₂” and the lowest energy “EuAu₅In” models. A comparison of the DOS curves for both models reveals significant contributions from the Au- d atomic orbitals (AOs), which mainly reside between -6.30 and -2.18 eV for “EuAu₄In₂” and between -6.57 and -2.01 eV for “EuAu₅In”. The states near the Fermi level, E_F , originate primarily from the Au- d states with minor contributions from the Eu- d and In- p AOs, respectively. An analysis of the VASP-based DOS curves for a hypothetical model of “Eu₅Au₁₆In₆”, which is derived from the orthorhombic structure of **III** through assignments of the major component indium on the disordered M sites (Wyckoff positions $8f$, $8f$, $8f$), reveals alike tendencies (Figure S10, [Supporting Information](#)). In particular, the states below E_F arise mainly from the Au- d AOs (between -2.33 and 2.32 eV), whereas the states around the Fermi level stem mostly from the bands accounting for the Au- d AOs with minor shares from the Eu- d and In- p states, respectively. Because the refined composition of **III** has 41.4 valence electrons per formula unit (VEs/fu) and the composition of “Eu₅Au₁₆In₆” corresponds to 44 VEs/fu, it is rather inapt to evaluate stability tendencies for the Eu₅Au₁₆(Au/In)₆ structures based on the relative position of the Fermi level in “Eu₅Au₁₆In₆”; however, E_F in “Eu₅Au₁₆In₆” falls in a deep pseudogap, indicating an electronically favorable situation.

Notwithstanding that the general features are similar for “EuAu₄In₂” and “EuAu₅In”, still, the models differ in both composition and valence electron counts (vec). The Fermi level in “EuAu₄In₂” (12 valence electrons) falls close to a local maximum in the DOS curve, while E_F in “EuAu₅In” (10 valence electrons) falls in a pseudogap that is generated in the DOS regions for vec of 9.9–10.6 VEs/fu. As the tetragonal structure

Table 3. Distance and –ICOHP/Bond Ranges, Average –ICOHP/Bond, Cumulative –ICOHP/Cell, and Percentage Contributions of the Diverse Interactions in “EuAu₄In₂”, “EuAu₅In-1”, “EuAu₅In-2”, “EuAu₅In-3”, and “EuAu₅In-4”^a

interaction	distance range [Å]	–ICOHP/bond range [eV/bond]	ave. –ICOHP/bond [eV/bond]	cum. –ICOHP	percentage
“EuAu ₄ In ₂ ”					
Au–Au	2.853–3.071	0.8351–0.7765	0.8058	12.8928	32.82
Au–Eu	3.187	0.3602	0.3602	2.8816	7.34
Au–In	3.021	0.6327	0.6327	20.2464	51.54
In–In	2.840	0.8155	0.8155	3.2620	8.30
“EuAu ₅ In” model 1					
Au–Au	2.751–3.012	1.0496–0.4811	0.7586	27.3091	63.19
Au–Eu	3.103–3.194	0.3747–0.1915	0.2857	1.7144	3.97
Au–In	2.796–2.986	1.0148–0.6927	0.8317	13.3076	30.79
Eu–In	3.159	0.4418	0.4418	0.8836	2.05
“EuAu ₅ In” model 2					
Au–Au	2.702–3.093	1.6429–0.5955	0.8740	32.3369	69.63
Au–Eu	2.988–3.200	0.5571–0.3650	0.4856	2.9134	6.27
Au–In	2.790–3.006	1.2100–0.5025	0.7503	10.5046	22.62
Eu–In	3.282	0.1580	0.1580	0.3160	0.68
In–In	3.117	0.3700	0.3700	0.3700	0.80
“EuAu ₅ In” model 3					
Au–Au	2.844–3.013	1.0112–0.6365	0.7893	50.5130	57.05
Au–Eu	3.106–3.110	0.4051–0.3121	0.3600	5.7604	6.50
Au–In	2.794–2.983	0.9107–0.6484	0.8068	32.2714	36.45
“EuAu ₅ In” model 4					
Au–Au	2.770–3.010	0.9405–0.5833	0.7436	49.0786	60.55
Au–Eu	3.111–3.125	0.3688–0.3113	0.3378	5.4044	6.67
Au–In	2.741–3.003	0.9884–0.6159	0.7062	25.4232	31.37
In–In	2.945	0.5740	0.5740	1.1480	1.41

^aTables compiling all distances, –ICOHP/bond, and their respective multiplicities are provided in the Tables S3–S6, [Supporting Information](#).

II refines as EuAu_{4.75}In_{1.25(2)}, which corresponds to 10.5 VEs/fu, the Fermi level will be located near the higher part of the pseudogap in the “EuAu₅In” DOS. Likewise, E_F in EuAu_{4.93}In_{1.07(2)} (10.1 VEs/fu) and EuAu_{5.03}In_{0.97(2)} (9.9 VEs/fu) will occur near higher and lower parts of the pseudogap, respectively. This result alludes to a higher relevance of valence electron concentrations toward distributing gold and indium sites.

In summary, the total energy calculations point to a site preference of the 4d sites for the Au/In partitioning in the tetragonal structures. On the basis of an analysis of the DOS curves for both “EuAu₄In₂” and “EuAu₅In” valence electron concentrations appear to play a relevant role in these structures. Another important contribution that one would expect to influence the assignments of elements on intrinsic positions within one structure is the bond energy.⁷³ To analyze its impact on the tetragonal structure in more detail, we followed up with a (chemical) bonding analysis on the “EuAu₄In₂” and “EuAu₅In-3” models.

Bonding Analysis. Analysis of the –COHP curves (Figure 4) and their integrated values (Table 3) for both “EuAu₄In₂” and “EuAu₅In” model-3 indicates that significant bonding interactions are evident for the homoatomic M–M as well as the heteroatomic Au–In contacts. A direct comparison between the –ICOHP values of both systems cannot be made as the average electrostatic potential in each DFT-based calculations is scaled to an arbitrary “zero” energy;⁷⁴ however, projecting –ICOHP values weighted by bond frequencies as percentages to the total bonding capacities has been demonstrated to provide conclusive hints to the bonding differences between dissimilar structures.^{53,54,75}

The Au–Au –ICOHP values range from 0.8351 to 0.7765 eV/bond for “EuAu₄In₂” and from 1.0112 to 0.6365 eV/bond for the more Au-rich composition and contribute 32.82% and 57.05% to the total bonding of the respective structures. Note that the percentage contribution of the Au–Au contacts is larger in “EuAu₅In” than in the more In-rich structure as a consequence of the higher Au–Au bond frequencies in “EuAu₅In”. Because –ICOHP values tend to scale similarly to bond strength, generally, the magnitude of the –ICOHP values will decrease as the bond lengths increase. For example, the largest Au–Au –ICOHP value of the ordered “EuAu₄In₂” (0.8351 eV/bond) occurs for the separations ($d = 2.853$ Å) that enclose the squared bases of the Au₈ cages within the *ab* plane, while the Au–Au separations in the zigzag chains between these Au₄ squares are 3.071 Å and have an –ICOHP value of 0.7765 eV/bond. The homoatomic Au–Au interactions are optimized at –1.18 eV for “EuAu₄In₂” and at –1.58 eV for “EuAu₅In” and remain nonbonding until E_F . The strongly antibonding Au–Au interactions at around –3.26 eV for “EuAu₄In₂” and –3.05 eV for “EuAu₅In” stem from the repulsions of the Au-d orbitals, which are observed typically for gold-rich intermetallic compounds.⁷⁶

The heteroatomic Au–In contacts are nonbonding around E_F (between –0.90 and 0.39 eV for “EuAu₄In₂” and between –0.44 and 1.55 eV for “EuAu₅In”) and change to antibonding states at 0.39 eV for “EuAu₄In₂” and at 1.55 eV for “EuAu₅In”. These interactions contribute 51.54% and 36.45% to the total bonding of “EuAu₄In₂” and “EuAu₅In-3”, respectively. The difference in the percentage contributions of the Au–In contacts to the total bonding capacities originates from the dissimilar bond frequencies between “EuAu₅In” and “EuAu₄In₂”, which has twice as many Au–In as Au–Au contacts.

Another significant difference between the “EuAu₄In₂” and the “EuAu₅In” models is the absence of any close, homoatomic In–In contacts between the face-sharing M@Au₈ cages in the latter structure. These contacts ($d = 2.840 \text{ \AA}$) are evident for the ordered “EuAu₄In₂” and show bonding interactions with an –ICOHP value of 0.8155 eV/bond; in the more gold-rich model, however, the short Au–In separations (2.794 Å) exhibit strong bonding populations between the M@Au₈ cages.

A comparison of the Au–Au and Au–In –ICOHP values for the “EuAu₄In₂” and the lowest energy “EuAu₅In” models reveals that these values scale in the same range for the more Au-rich structure but show a significant difference in the ordered “EuAu₄In₂” (Table 3). As a consequence, there are strong Au–Au as well as Au–In bonding interactions in “EuAu₅In”, whereas the –ICOHP values are much smaller for the heteroatomic than for the homoatomic contacts in the ordered “EuAu₄In₂”. Although bond frequencies are a substantial factor in these structures and for their formations, this result also indicates the influence of the local symmetry and the intrinsic atomic arrangements. This impact comes even clearer from a comparison between “EuAu₅In” models 1 and 2. Even though model 1 has a higher Au–In bond frequency than model 2, this discrepancy is overcome by the higher amount of strong Au–Au interactions in 2, for which a lower total energy is evaluated rather than for model 1. This outcome suggests that the subtle interplay between Au–In bond frequencies and local atomic environments not only adjusts the site preferences for the Au/In partitioning but also controls the structural preferences for these compounds. The formally anionic networks of the diverse models encapsulate europium atoms, which reside in the centers of two face-sharing M₈ cages. The magnitudes of the M–Eu –ICOHP values are smaller than those for the Au–Au, Au–In, and In–In interactions, which is observed typically for rare-earth late-transition-metal intermetallic compounds.⁶⁸

CONCLUSIONS

Exploration of the gold-rich regions for the Eu–Au–In system lead to the discovery of three ternary intermetallic compounds. EuAu_{0.46}In_{1.54(2)} (I) can be derived from the binary EuAu₂ through partial indium occupation of the Au sites. Significant gold distributions on the indium sites are observed for the “EuAu₄In₂”-derived EuAu_{4+*x*}In_{2–*x*} with $x = 0.75(2)$ (II), 0.93(2), and 1.03(2). The mutual exchange of Au and In atoms for these structures while maintaining their host lattices substantiates the mutual exchangeabilities of Au and In but also points to certain electronic flexibilities for I and II. A combination of the structural elements of I and II was realized in Eu₅Au_{17.29}In_{4.71(3)} (III), which adopts its own structure type. A topological analysis for the recently reported Eu₅Au_{17.7}In_{4.3} and Eu₅Au_{17.29}In_{4.71(3)} (III) reveals that these compounds formally represent one-dimensional intergrowths of the series EuAu₂–“EuAu_{4+*x*}In_{2–*x*}”.

The electronic structure was examined for the EuAu_{4+*x*}In_{2–*x*} compounds, which show the shortest repeat unit along (001) within the reported series. Electronic structure calculations on diverse “EuAu₅In” models indicated that there are clear site preferences of mixing indium on the 4*d* site rather than the 8*h* sites. From a topological analysis it becomes clear that this tendency arises from the attempt to optimize the frequency of the heteroatomic Au–In bonds, which is achieved through partial occupation of the 4*d* sites by indium. The positions of the Fermi level in the DOS of the “EuAu₄In₂” and the lowest

energy “EuAu₅In” models reveal electronically favorable situations for the tetragonal structures but also suggest that valence electron concentrations are of higher relevance for their formations. On the basis of a chemical bonding analysis it is the subtle interplay between the heteroatomic bond frequencies and the approach to sustain strong local atomic contacts that ascertains the adaptation of an intrinsic structure for this system.

ASSOCIATED CONTENT

Supporting Information

Tables of the atomic anisotropic displacement parameters for I, II, and III; simulated and measured PXRD patterns of I, II, and III; representations of the miscellaneous coordination environments in II and III; spin-polarized (VASP) and TB-LMTO-based DOS data for “EuAu₄In₂” and the lowest energy “EuAu₅In” model; distances, multiplicities and –ICOHP/bond values for selected interactions in “EuAu₄In₂” and all “EuAu₅In” models; DOS and projected DOS curves (VASP) of a hypothetical composition “Eu₅Au₁₆In₆”; crystallographic information on the ternary EuAu_{0.46}In_{1.54(2)} (I), EuAu_{4+*x*}In_{2–*x*} ($x = 0.75(2)$ (II), 0.93(2), 1.03(2)) and Eu₅Au_{17.29}In_{4.71(3)} (III) in cif form; The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.inorgchem.5b00257.

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Notes

The authors declare no competing financial interest.

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